# Exploring new physics with pulsar timing arrays.

"Quarkonia meet Dark Matter" workshop, TUM

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Based on work with Torsten Bringmann, Paul Frederik Depta, Thomas Konstandin, Kai Schmidt-Hoberg and Pedro Schwaller

JCAP 11 (2023) 053 and [2306.17836]

March 21, 2024



#### Outline of this talk.

- 1. The PTA signal
- 2. Phase transitions vs. precision cosmology
- 3. Clustered PBHs
- 4. BSM or boring?



## In case you haven't heard the news.



### Pulsar timing arrays.

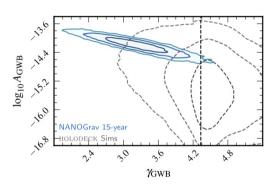


Millisecond pulsars emit radio pulses with an extremely stable frequency

- GWs affect propagation time → modulate observed pulse frequency
- PTAs monitor pulse frequency using radio telescopes on Earth
  - Fit pulse data with timing model
- Fourier decomposition of timing residuals shows common spectrum, which is due to GWs

## GW background from supermassive black hole binaries.

The observed GW spectrum is consistent with a power-law of amplitude A and slope  $\gamma$ 

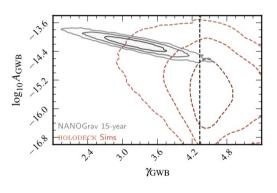


[NANOGrav collaboration, 2023]

## **GW** background from supermassive black hole binaries

- The observed GW spectrum is consistent with a power-law of amplitude A and slope  $\gamma$
- But: Astrophysical simulations based on realistic BH populations predict much weaker signals with higher γ

The standard explanation is off! Other signal sources?

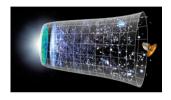


[NANOGrav collaboration, 2023]

#### Possible non-standard sources of the nHz background.

#### Inflation

Reentering of tensor fluctuations

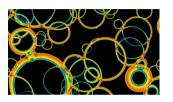


**Topological defects**Cosmic strings and domain walls



**Phase transitions** 

Connection to dark matter?



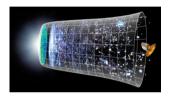
**Primordial black holes**But only if they are clustered...



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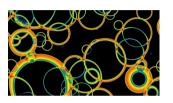


**Topological defects**Cosmic strings and domain walls



#### **Phase transitions**

Connection to dark matter?



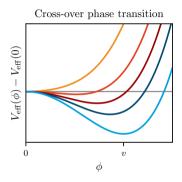
#### **Primordial black holes**

But only if they are clustered...

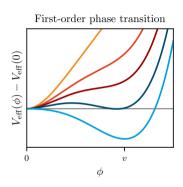




### Cosmological phase transitions.



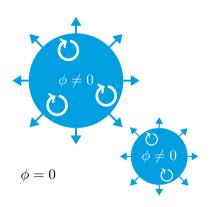
A scalar field "rolls down" from  $\phi=0$  to  $\phi=v$ , when the bath cools from high temperatures to low temperatures.

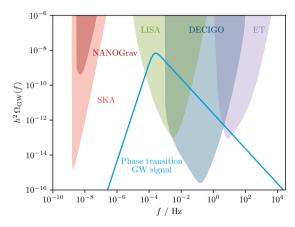


A scalar field tunnels to the true potential minimum ( $\phi \neq 0$ ) to minimize its action ( $\sim$  free energy).

#### **Gravitational waves from first-order phase transitions.**

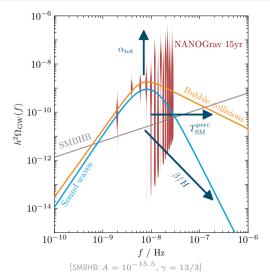
Bubbles of the new phase nucleate, collide and perturb the plasma...





... giving rise to a stochastic gravitational wave background which can be observed.

#### Parametrization of the GW signal.

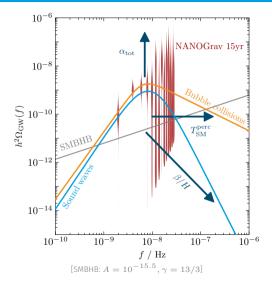


$$h^2\Omega_{\rm GW}^{\rm SW,bW}(f) \simeq 10^{-6} \left(\frac{\alpha}{\alpha+1}\right)^2 \left(\frac{H}{\beta}\right)^{1,2} \mathcal{S}\left(\frac{f}{f_{\rm peak}}\right)$$
 with  $f_{\rm peak} \simeq 0.1\,{\rm nHz} \times \frac{\beta}{H} \times \frac{T}{{\rm MeV}}$ 

To fit the new pulsar timing data:

- Strong transitions,  $lpha \simeq rac{\Delta V}{
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- Slow transitions,  $\beta/H \approx 10$
- Percolation around  $T \approx 10\,\mathrm{MeV}$

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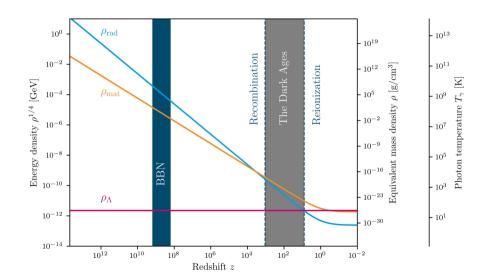
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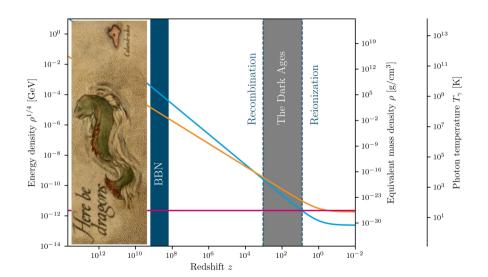
But there's no SM phase transition at 10 MeV?!



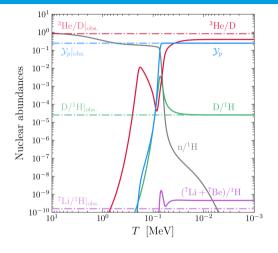
#### A brief history of time.



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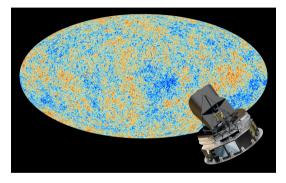
### The Big Bang Nucleosynthesis and the CMB.



[Paul Frederik Depta, 2021]

- Observations of primordial light element abundances in good agreement with standard BBN
- $N_{
  m eff}^{
  m BBN} = 2.898 \pm 0.141$  [Yeh+, 2207.13133]

### The Big Bang Nucleosynthesis and the CMB.



[ESA and the Planck Collaboration, D. Ducros]

- Observations of primordial light element abundances in good agreement with standard BBN
- $N_{
  m eff}^{
  m BBN} = 2.898 \pm 0.141$  [Yeh+, 2207.13133]
- $\cdot~N_{
  m eff}^{
  m CMB} = 2.99 \pm 0.17$  [Planck, 1807.06209]
- Consistent with  $N_{
  m eff}^{
  m SM}=3.044$  from 3 u generations [Bennet\*, 2012.02726]

Thermalized BSM species at  $T \lesssim 1\,\mathrm{MeV}$  are ruled out. Before that: no constraints.

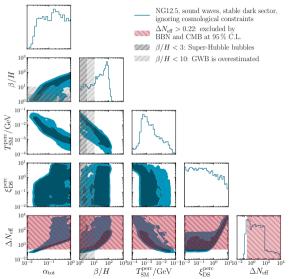
### Let's put the transition in a dark sector.

- · SM has no MeV phase transition  $\rightsquigarrow$  Assume a weakly coupled  $\mathcal{O}(\text{MeV})$  scalar!
- $\cdot$  Dark sector temperature is crucial for GW prediction,  $T_{
  m DS}=\xi_{
  m DS}~T_{
  m SM}$  [CT+, 2109.06208]
- Stable dark sector: additional DS energy density accelerates expansion and changes early element abundances and CMB anisotropies through

$$\Delta N_{
m eff} pprox 6 imes \left( lpha + rac{1+lpha}{10} \xi_{
m DS}^4 
ight) \; , \quad \Delta N_{
m eff} < 0.22 \; @95 \, \% \; {
m C.L.}$$

• Decaying dark sector: Energy transfer to the SM plasma, changing element abundances and CMB anisotropies. Constraints require  $au < 0.1\,\mathrm{s.}$  [Depta+, 2011.06519]

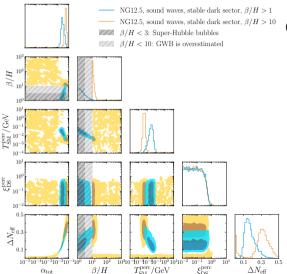
#### The tension between PTAs, CMB and BBN.



- Performed fit of the pulsar data with NANOGrav's own code enterprise
- $\red{\uparrow}$  A good fit requires an enormous reheating of the dark sector:  $\Delta N_{\rm eff}$  can grow arbitrarily large
- $\raiseta$  Bubble sizes would need to be super-Hubble to be okay with  $\Delta N_{
  m eff}$  Causality  $\raiseta$  GW prediction  $\raiseta$

The tension cries for a global fit

#### Global fits.

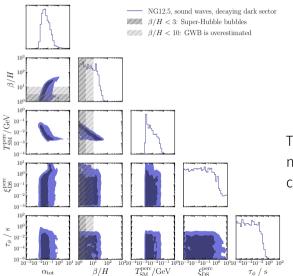


Global fit = compute global maximum of

$$\begin{split} \mathcal{L}_{\text{glob}}(\vec{\theta}_{\text{PSR}}, \vec{\theta}_{\text{PT}}) &= \\ \mathcal{L}_{\text{PTA}}(\vec{\theta}_{\text{PSR}}, \vec{\theta}_{\text{PT}}) \times \mathcal{L}_{\text{cosmo}}(\Delta N_{\text{eff}}(\vec{\theta}_{\text{PT}})) \end{split}$$

- $\beta/H > 1$ : would be a good fit, if the GW spectrum were reliable
- $\beta/H > 10$ : spectra reliable, but not having a phase transition is better than violating BBN and CMB bounds!

### Decays to the rescue.

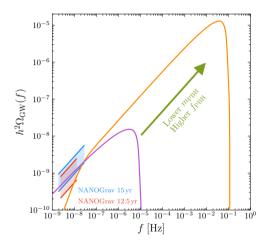


## Decays save the fit!

They only need to happen before neutrino decoupling,  $T_{\rm SM}\gtrsim 2\,{\rm MeV}$ , corresponding to fast decays,  $\tau\lesssim 0.1\,{\rm s}$ .



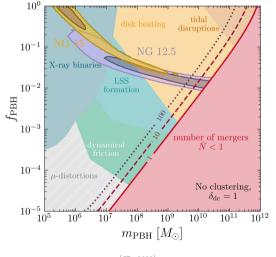
### Gravitational waves from primordial black hole mergers.



- Inflation leaves large super-Hubble density perturbations
- Black holes form when these come into causal contact again, long before the death of the first stars
- Described by (monochromatic) mass m<sub>PBH</sub> and DM fraction f<sub>PBH</sub>



### PBHs without clustering cannot explain the PTA data.



- Scan over  $m_{PBH}$  and  $f_{PBH}$
- Region favored by PTAs is excluded by astrophysical bounds
- Crucial: exclude regions with small merger numbers. (Atal et al. came to the wrong conclusion [2012.14721])

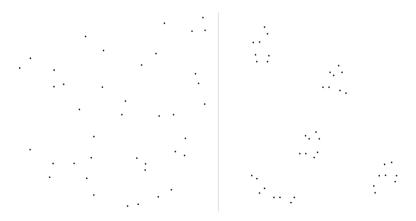
Homogeneously distributed PBHs cannot explain the PTA data!

2023]

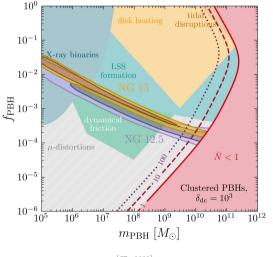
### What is clustering?

$$\delta_{
m dc}=1$$
: Poisson-distributed PBHs

$$\delta_{ extsf{dc}} = 1 + rac{\delta n_{ extsf{PBH}}^{ ext{loc}}}{ar{n}_{ extsf{PBH}}} \gg 1$$
: Clustering



## Clustered PBHs can explain the PTA data.



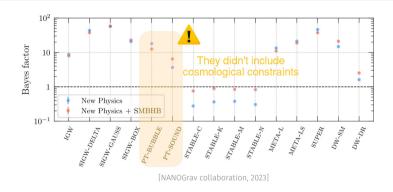
- Clustering increases the merger rates, requiring less PBHs to explain the signal: shift to smaller  $f_{\rm PBH}$
- Astrophysical bounds are dubious for clustering
- Fermilab group pointed out that  $\mu$ -distortion bounds are model dependent [Hooper+, 2308.00756]

Clustered PBHs can explain the PTA data!

+, 2023]



#### The evidence for new physics.



- New physics matches spectra better than standard SMBHBs
- We should perform global fits, including additional constraints & astrophysical parameters

Still: As soon as a single merger or strong anisotropy is found in the data, all cosmological explanations will be practically dead.

## Take-home messages.

- New physics can explain the signal better than astrophysics
- Stable dark sector phase transition explanations for PTA data are in tension with precision cosmology
- Decaying dark sectors with  $\tau < 0.1\,\mathrm{s}$  are a viable option, can compete with SMBHBs
- Primordial black hole mergers can only explain the signal if they are clustered

Thank you very much for your attention!

Do you have any questions?





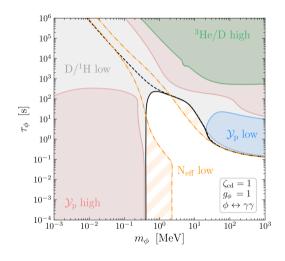
## How the density contrast increases the merger rate

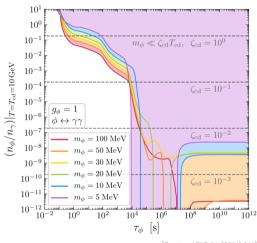
$$\begin{split} \Omega_{\text{GW}}(f) &= \frac{f}{\rho_{\text{crit}}} \int_{0}^{t_0} \mathrm{d}t \left[ R(t) \left. \frac{\mathrm{d}E_{\text{GW}}}{\mathrm{d}f_{\text{r}}} \right] \right|_{f_r = (1+z)f} \\ R(t) &= \int_{0}^{\tilde{x}} \mathrm{d}x \int_{x}^{\infty} \mathrm{d}y \frac{\partial^{2}n_{3}}{\partial x \, \partial y} \delta(t - \tau(x, y)) \\ &\propto \frac{\delta_{\text{dc}}^{16/37}}{\tilde{x}^{3}\tilde{\tau}} \left( \frac{t}{\tilde{\tau}} \right)^{-34/37} \left( \Gamma \left[ \frac{58}{37}, \frac{4\pi}{3} \tilde{x}^{3} \delta_{\text{dc}} n_{\text{PBH}} \left( \frac{t}{\tilde{\tau}} \right)^{3/16} \right] - \\ &\Gamma \left[ \frac{58}{37}, \frac{4\pi}{3} \tilde{x}^{3} \delta_{\text{dc}} n_{\text{PBH}} \left( \frac{t}{\tilde{\tau}} \right)^{-1/7} \right] \right) \end{split}$$

#### With:

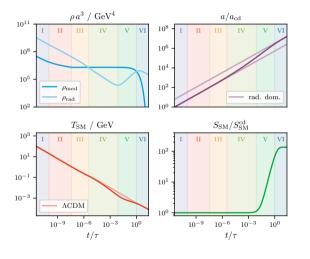
- $\cdot$   $\delta_{
  m dc} \simeq rac{n_{
  m PBH}^{
  m loc}}{ar{n}_{
  m PBH}^{
  m loc}}$ : Density contrast
- $\cdot \ x, (y)$ : comoving distance of (next-to-) nearest neighbor PBH
- $\tilde{x}$ : farthest comoving distance two PBHs can have
- $\tilde{\tau}$ : Merger timescale

### Electromagnetic scalar decays at MeV temperatures.





#### The out-of-equilibrium decay of a dark mediator.

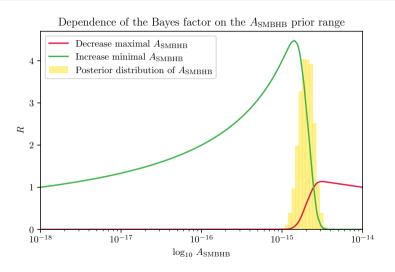


Energy densities  $\rho_i(t) \stackrel{\text{sets}}{\leadsto} \text{Scale factor}$   $a(t) \stackrel{\text{sets}}{\leadsto} \text{Temperatures } T_{\text{SM/DS}}(t) \stackrel{\text{set}}{\leadsto}$  Particle content  $\stackrel{\text{sets}}{\leadsto} \rho_i(t) \stackrel{\text{sets}}{\leadsto} \dots$ 

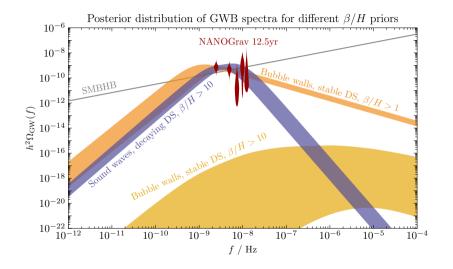
#### Six phases:

- I Relativistic mediator
- II Cannibalistic mediator
- III Non-relativistic mediator
- IV Early matter domination
- V Entropy injection
- VI Mediator decay

### How the choice of priors changes a Bayes factor.



### Why violins shouldn't be used for fits including cosmological constraints.



#### We can derive new bounds on clustered PBHs.

